

Hadroquarkonium from Lattice QCD

Francesco Knechtli

with Maurizio Alberti, Gunnar Bali, Sara Collins,
Graham Moir and Wolfgang Söldner

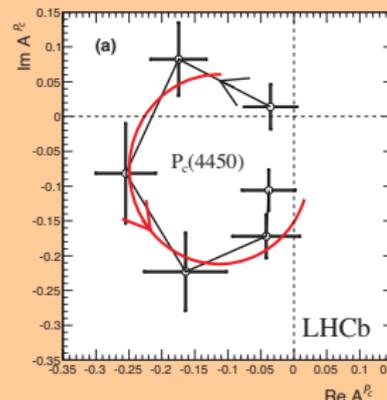
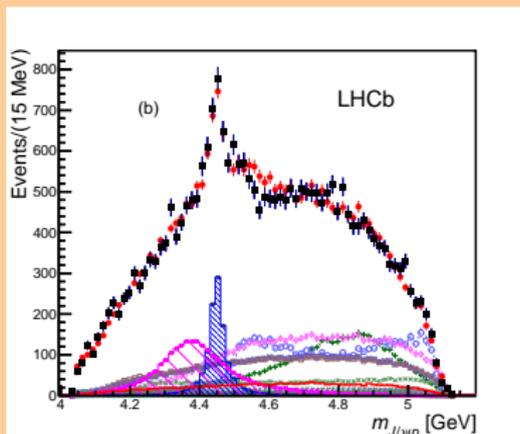
Wuppertal – Regensburg – Cambridge

QWG 2016, PNNL, Richland, June 8, 2016



Motivation

LHCb pentaquark candidates



$P_c^+(4380)$ ($J^P = \frac{3}{2}^-$) and $P_c^+(4450)$ ($J^P = \frac{5}{2}^+$) from $\Lambda_b \rightarrow J/\psi p K$ [LHCb: R. Aaij et al, 1507.03414]

Conjecture of attractive forces between charmonium and pp systems [Brodsky, Schmidt and de Teramond, PRL64 (90) 1011]

Many interpretations

Motivation

Hadro-quarkonia

5 ($4 q, 1 \bar{q}$) quark systems are very difficult to study directly on the lattice. Here we test a particular model instead.
Hadro-quarkonia: quarkonia bound “within” ordinary hadrons [S. Dubynskiy and M. Voloshin, 0803.2224]

Many combinations of baryon plus charmonium are close-by. Examples:

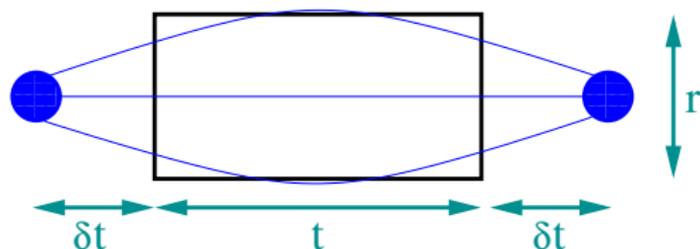
$J^P = \frac{3}{2}^-$: $m(\Delta) + m(J/\psi) \approx 4329 \text{ MeV vs. } 4380 \text{ MeV}$ (width 200 MeV)

$J^P = \frac{5}{2}^+$: $m(N) + m(\chi_{c2}) \approx 4496 \text{ MeV vs. } 4450 \text{ MeV}$ (width 40 MeV)



Correlation function

The hadro-quarkonium model can be tested in the static limit. To leading order in (p)NRQCD, quarkonia can be approximated by the non-relativistic Schrödinger equation with a static potential $V_0(r)$. Does this become more or less attractive, when light hadrons are “added”?



Create a zero-momentum projected hadronic state $|H\rangle$ at the time 0 .

Let it propagate to δt , create a quark-antiquark “string”.

Destroy this at $t + \delta t$ and the light hadron at $t + 2\delta t$:



Correlation function

$\bar{Q}Q$ binding energy shift “within” a hadron H

We compute

$$C(t; \delta t) = \frac{\langle H^\dagger(t + 2\delta t, \underline{x}) W(t, r) \bar{H}(0, \underline{0}) \rangle}{\langle H^\dagger(t + 2\delta t, \underline{x}) \bar{H}(0, \underline{0}) \rangle \langle W(t, r) \rangle}$$

where we average over the spatial positions of the Wilson loop $W(r, t)$ and over the hadronic sink positions \underline{x} .

The shift of the binding potential is obtained from

$$\Delta V_H(r, \delta t) = V_H(r, \delta t) - V_0(r) = - \lim_{t \rightarrow \infty} \frac{d}{dt} \ln C(t; \delta t).$$

and extrapolating $\delta t \rightarrow \infty$.



Lattice details

We analyse the $N_f = 2 + 1$ **CLS** ensemble C101 (96×48^3 sites) [Bruno et al, 1411.3982]:

$m_\pi = 220$ MeV, $m_K = 470$ MeV, $Lm_\pi \approx 4.6$, $L \approx 4.1$ fm,
 $t_0/a^2 = 2.9085(51)$, $a \approx 0.086$ fm (from t_0/a^2 extrapolated to physical point)

High statistics: 1552 configs, separated by 4 MDUs, times 12 hadron sources (1 forward, 1 backward, 10 forward and backward propagating \Rightarrow 22 2-point functions).

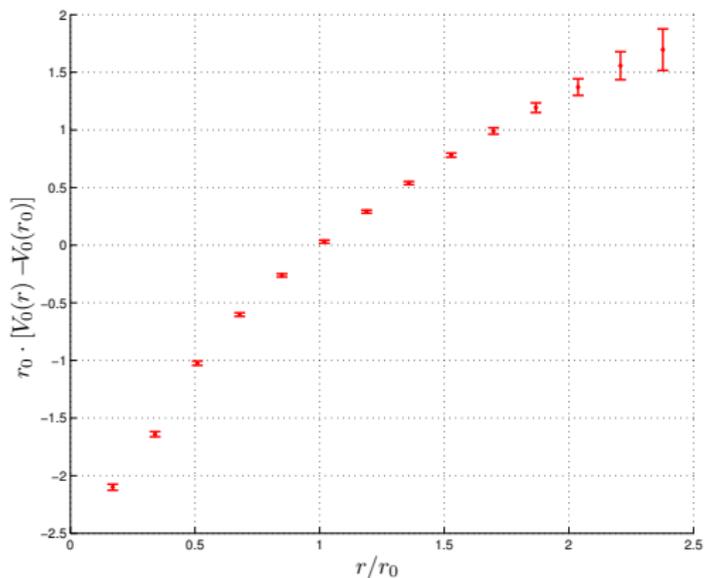
Wilson loops at all positions and in all directions.

Hadronic two-point functions have optimised overlap with the ground state.



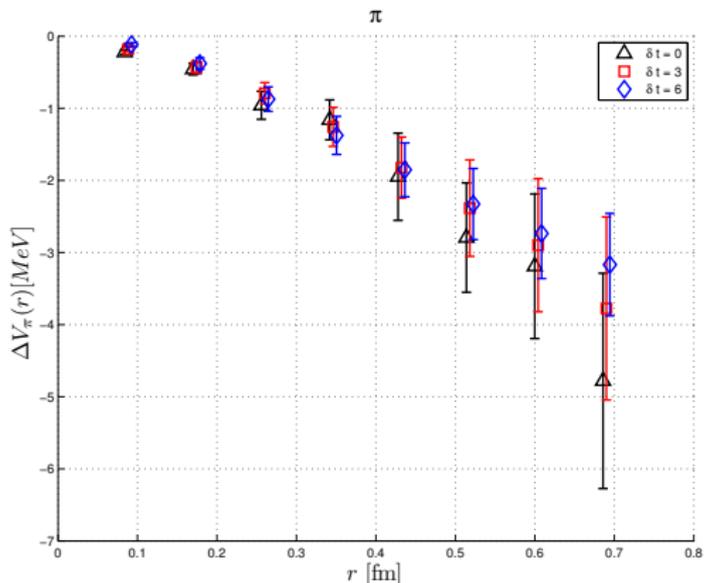
Static potential $V_0(r)$

Using the methods of [Donnellan, FK, Leder and Sommer 1012.3037] variational basis with 0, 5, 7, 12 spatial HYP levels; open boundary conditions: average between $t = 24$ and $t = 72$; Sommer scale $r_0/a = 5.890(41)$



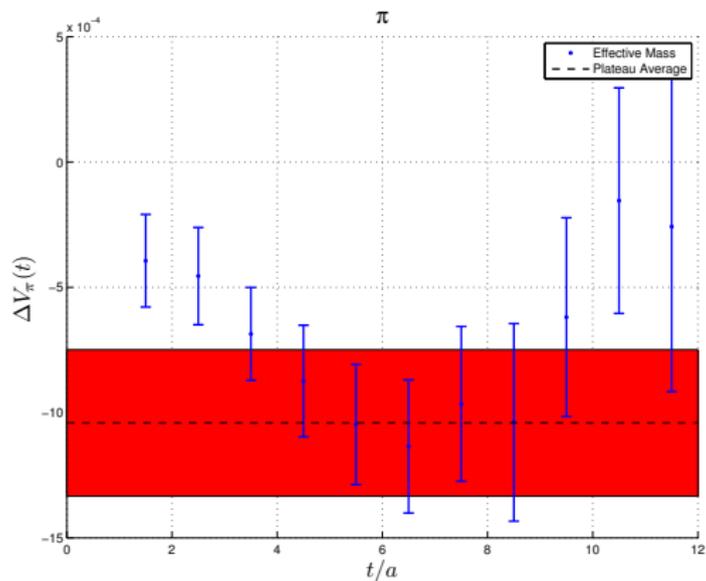
$\bar{Q}Q$ binding energy shift “within” a pion

[M. Alberti et al, in preparation]



We can resolve small energy differences

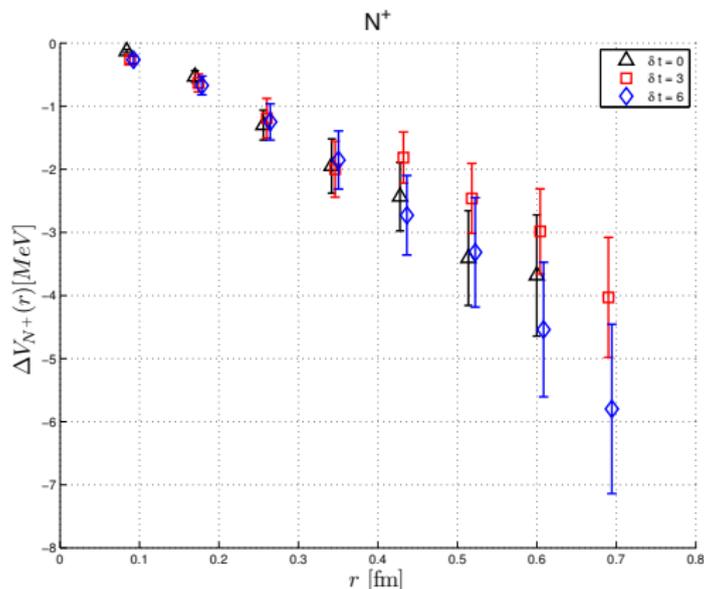
Details of the pion analysis



Example of effective masses at $\delta t = 3a$ and $r = 6a \approx r_0$



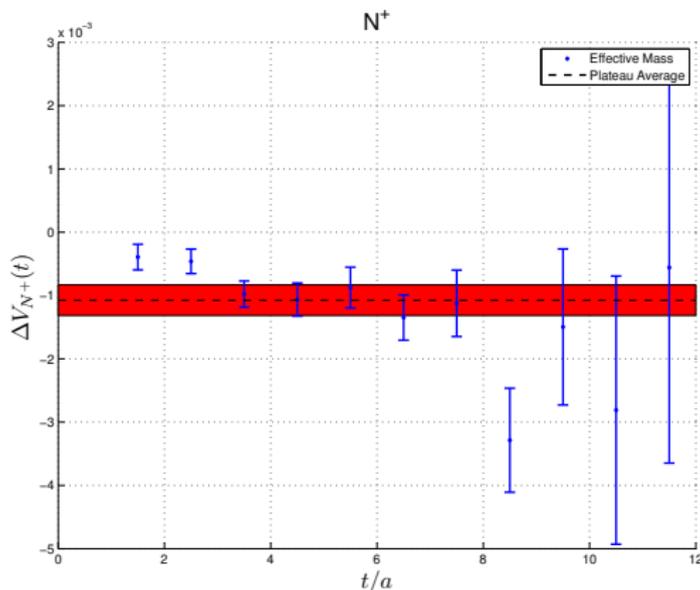
$\bar{Q}Q$ binding energy shift “within” a nucleon



For a baryon of parity P , parity P is taken in the forward propagator and parity $-P$ in the backward propagator



Details of the nucleon analysis



Example of effective masses at $\delta t = 3a$ and $r = 6a \approx r_0$



The baryon decuplet

Candidates for pentaquark states

We measure decuplet baryons with **helicity** $\pm \frac{3}{2}$.

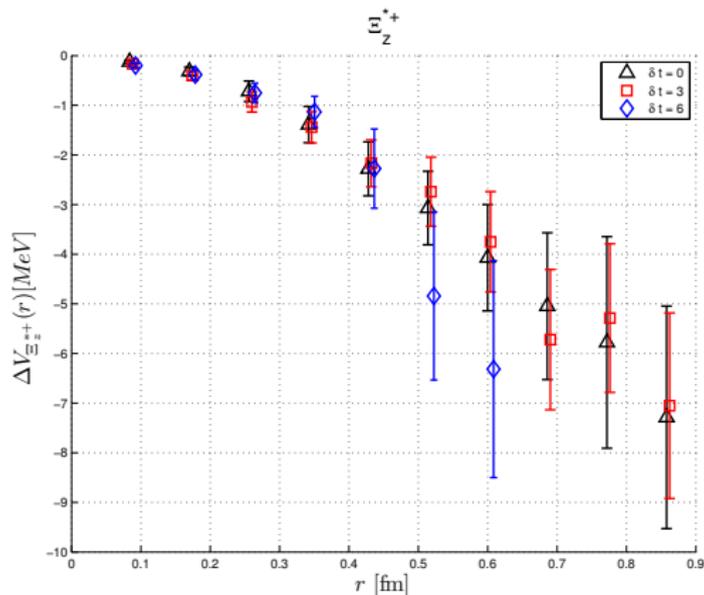
Quarkonium in S-wave has $J^P = 0^-$ and 1^- .

Combining 0^- or 1^- quarkonium with a $\frac{3}{2}^+$ decuplet baryon gives a $J^P = \frac{3}{2}^-$ state.

Combining 1^- quarkonium with a $\frac{3}{2}^-$ decuplet baryon gives a $J^P = \frac{5}{2}^+$ state.

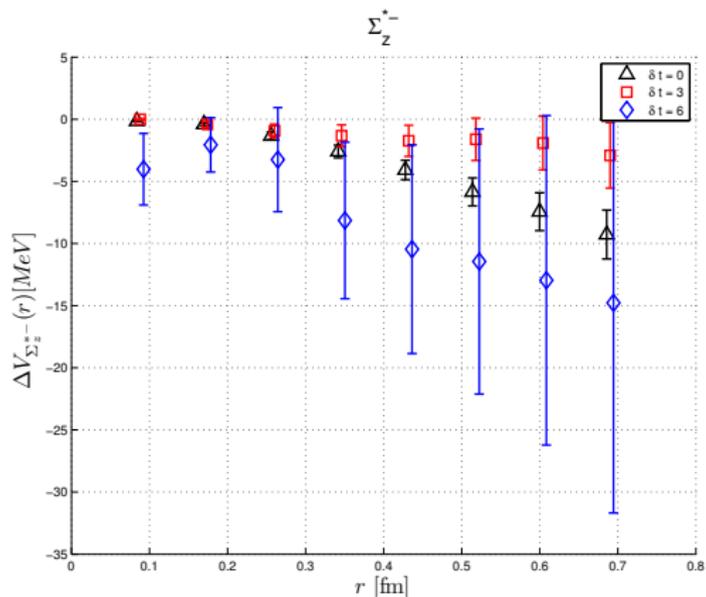


$\bar{Q}Q$ binding energy shift “within” a $\Xi^*(\frac{3}{2}^+)$



Polarisation: correlation of baryon polarised in z direction with Wilson loops in z direction

$\bar{Q}Q$ binding energy shift “within” a $\Sigma^*(\frac{3}{2}^-)$



The signal for $\frac{3}{2}^-$ decuplet baryons is noisier. Also it does not match the mass of the $\frac{5}{2}^+$ pentaquark.



Summary

Conclusions

- ▶ Test of the idea of hadro-quarkonium using lattice QCD in the static limit for the heavy quarks.
- ▶ Modifications of the static potential in the presence of light hadrons appear tiny $\approx -3 \text{ MeV}/(0.5 \text{ fm})$. Compared to the “slope” of the static potential $\approx 1 \text{ GeV}/\text{fm}$ it is a 0.6% reduction.
- ▶ Virial theorem for a linear potential \Rightarrow energy of states decreases by $|\Delta E| = 1-2 \text{ MeV}$. Somewhat inconsistent with the original hadro-charmonium.
- ▶ Interestingly, there is a similar attraction in all of the channels investigated so far.

Outlook

study of finite volume effects \longrightarrow Lattice 2016

